# COMPARATIVE MODAL ANALYSIS OF A RIGID STRUCTURE WITH CONSERVATIVE INSULATION

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# ABSTRACT

This paper deals with dynamics of compact rigid structures with elastic insulation and subjected to various vibratory, shock or seismic actions. It was supposed a simple rigid model with six degrees of freedom, affected by one or both vertical symmetry planes. Modal analysis was performed for six different rigidity cases. Concluding remarks dignify the correlation between isolation characteristics and natural pulsation in terms of eigenvalues and eigenvectors.

KEYWORDS: modal analysis, dynamic behaviour, rigid structure, eigenvalue, eigenvector

## **1. INTRODUCTION**

The structural systems are made of material systems as solid bodies (rigid or deformable) with elastic or viscoelastic linkages subjected to external dynamic actions, having as consequences the generation of inertial effects.

Dynamic behavior of structural systems is described by mathematical equations that take different forms from a specific case to another. The representation in calculations of the solid bodies (rigid or deformable), of the linkages and of the whole system is based on the concept of dynamic system and dynamic model.

The dynamic system is an abstraction of the physical and mechanical characteristics of the structural system whose mechanical condition changes during time.

Any dynamical system is characterized by some specific qualitative properties (inertial, dissipative, elastic) represented by the values of measurable parameters (mass, moments of inertia, the damping coefficient, the rigidity/ flexibility coefficient).

The dynamic model is essentially an idealized form, simplified or schematized of a dynamic system in order to reduce the numerical analysis operations without that the real processes (qualitative and quantitative) being significantly modified.

The dynamic response is the instantaneous state of a dynamic system over which have been applied external dynamic actions, real and variable during time. The dynamic response can be expressed through fundamental kinematic parameters (displacement, velocity, acceleration) or through derived parameters (energy, sectional strains, stresses, deformations, generalized forces).

#### **2. THEORETICAL APPROACHES**

It is proposed the complex dynamic model of a rigide structure with six degrees of freedom which consists of three translational linear coordinates x, y, z and three rotational angular coordinates  $\varphi_x, \varphi_y, \varphi_z$ . For the proposed model, we will study the behavior of the structure under the vibration action and in the presence of elastic elements.

Thus, the motion equations of the rigid with elastic linkages are written:

$$\underline{A}\underline{\ddot{q}} + \underline{C}\underline{q} = \underline{\underline{0}} \tag{1}$$

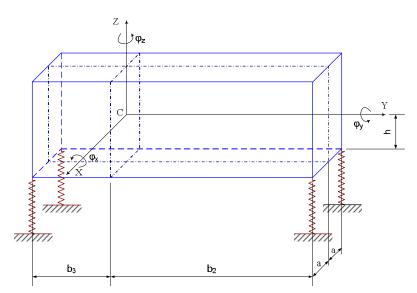


Fig.1. Model of the rigid system with six degrees of freedom, elastic supported in four points on inferior base, with a longitudinal vertical plane of symmetry

In an analytical form, the system is:

$$\begin{cases} m\ddot{x} + x\sum_{ix} k_{ix} + \varphi_{y}\sum_{ix} k_{ix}z_{i} - \varphi_{z}\sum_{ix} k_{ix}y_{i} = 0 \\ m\ddot{y} + y\sum_{iy} k_{iy} - \varphi_{x}\sum_{iy} k_{iy}z_{i} + \varphi_{z}\sum_{iy} k_{iy}x_{i} = 0 \\ m\ddot{z} + z\sum_{iz} k_{iz} + \varphi_{x}\sum_{iz} k_{iz}y_{i} - \varphi_{y}\sum_{iz} k_{iz}x_{i} = 0 \\ J_{x}\ddot{\varphi}_{x} - y\sum_{iy} k_{iy}z_{i} + z\sum_{iz} k_{iz}y_{i} + \varphi_{x}\sum_{iz} (k_{iy}z_{i}^{2} + k_{iz}y_{i}^{2}) - \varphi_{y}\sum_{iz} k_{iz}x_{i}y_{i} - \varphi_{z}\sum_{iy} k_{iy}z_{i}x_{i} = 0 \\ J_{y}\ddot{\varphi}_{y} + x\sum_{ix} k_{ix}z_{i} - z\sum_{iz} k_{iz}x_{i} - \varphi_{x}\sum_{iz} k_{iz}x_{i}y_{i} + \varphi_{y}\sum_{iz} (k_{iz}x_{i}^{2} + k_{iz}z_{i}^{2}) - \varphi_{z}\sum_{ix} k_{ix}y_{i}z_{i} = 0 \\ J_{z}\ddot{\varphi}_{z} - x\sum_{ix} k_{ix}y_{i} + y\sum_{iy} k_{iy}x_{i} - \varphi_{x}\sum_{iy} k_{iy}z_{i}x_{i} - \varphi_{y}\sum_{ix} k_{ix}y_{i}z_{i} + \varphi_{z}\sum_{ix} (k_{ix}y_{i}^{2} + k_{iy}x_{i}^{2}) = 0 \end{cases}$$

$$(2)$$

The system (2) is difficult to solve analytically or using the matriceal formalism because it requires a large amount of calculation, and the sixth degree polynomial equation of the natural pulsations involves difficulties in solving and analysis. The solution is the automatic numerical calculus of the differential motion equations system of second degree, resulting a system with 12 differential equations of first degree, which can be integrated without difficulty. On the other hand, at the use of numerical analysis appears as a disadvantage the highlighting of the influence of the dynamic system physical characteristics. Thus, the analysis is done by repeated tests, using different sets of values for the input data. To this end, both for the elimination of coupling movements and to analytically solve the dynamic system model, certain sized and structure requirements may be imposed to the system, leading to a decoupling of the equation system into simple subsystems, easier to integrate.

As discussed above, we consider the case

of the rigid structure, elastically supported in four points on inferior base, with a longitudinal vertical plane of symmetry yCz as in Figure 1. In this situation a few simplifying assumptions are valid:

- the dimensions of the analyzed rigid structure are symmetrical in relation to the considered plane
- the elastic linkages are identical, have symmetrical positions and are located in the same horizontal plane

Due to the mentioned symmetries, a part of coupling terms from stiffness matrix are canceled, and we have:

$$\sum k_{iy} x_i = 0$$

$$\sum k_{iz} x_i = 0$$

$$\sum k_{iz} x_i y_i = 0$$

$$\sum k_{iy} z_i x_i = 0$$
(3)

Through the disappearance of the coupling terms, the system decouples into two

subsystems described by coordinates (y, z,  $\varphi_x$ ) and (x,  $\varphi_y$ ,  $\varphi_z$ ). written the equations of the free vibrations. Thus, we have for the subsystem (y, z,  $\varphi_x$ ):

For the two decoupled subsystems can be

$$\begin{cases} m\ddot{y} + 4k_{y}y + 4hk_{y}\varphi_{x} = 0\\ m\ddot{z} + 4k_{z}z + 2k_{z}(b_{3} - b_{2})\varphi_{x} = 0\\ J_{x}\ddot{\varphi}_{x} + 4hk_{y}y + 2k_{z}(b_{3} - b_{2})z + 2[k_{z}(b_{2}^{2} + b_{3}^{2}) + 2h^{2}k_{y}]\varphi_{x} = 0 \end{cases}$$
(4)

and for the subsystem (x,  $\varphi_{y}$ ,  $\varphi_{z}$ ):

$$\begin{cases} m\ddot{x} + 4k_{x}x - 4hk_{x}\varphi_{y} - 2k_{x}(b_{3} - b_{2})\varphi_{z} = 0\\ J_{y}\ddot{\varphi}_{y} - 4hk_{x}x + 4(h^{2}k_{x} + a^{2}k_{z})\varphi_{y} + 2hk_{x}(b_{3} - b_{2})\varphi_{z} = 0\\ J_{z}\ddot{\varphi}_{z} - 2k_{x}(b_{3} - b_{2})x + 2hk_{x}(b_{3} - b_{2})\varphi_{y} + 2[2a^{2}k_{y} + k_{x}(b_{2}^{2} + b_{3}^{2})]\varphi_{z} = 0 \end{cases}$$
(5)

Further it is proposed an analyze of the vibrations of the two subsystems characterized each by three dynamic coordinates (degrees of freedom) coupled.

For each of the two subsystems, with elastic linkages and three degrees of freedom, the vector of the generalized coordinates is:

$$\underline{\underline{q}} = [q_1, q_2, q_3]^T \tag{6}$$

Using the classical mathematical apparatus, were written the quadratic forms of the system energies and then the case II Lagrange equations were used for obtaining motion equations, written matriceal in the form (1). The solution for the system (1) has been sought as:

$$\underline{\underline{q}} = \underline{\underline{a}} \sin pt = \begin{cases} a_1 \\ a_2 \\ a_3 \end{cases} \sin pt \tag{7}$$

where 
$$\underline{a} = \begin{cases} a_1 \\ a_2 \\ a_3 \end{cases}$$
 is the vector of the motion

amplitudes.

Taking into account the proposed form of the system solution, equation (1) becomes:

$$\left(\underline{C} - p^2 \underline{A}\right)\underline{\underline{a}} = \underline{\underline{0}}$$
(8)

Equation (8) has nonzero solutions only if the determinant of the matrix is zero. This equation is third degree with  $p^2$  as variable and represents the equation of the natural pulsations of the dynamic system with three degrees of freedom. Through analytically or numerically resolving the equation (8), we obtain the three natural pulsations of the system  $p_1$ ,  $p_2$ ,  $p_3$ .

#### **3. CASE STUDY**

The case study was made for two types of symmetry of the proposed structure, namely:

considering the structure having a longitudinal vertical plane of symmetry, case for which as numerical values were proposed:

a=7.5 m  $b_{3}=12 m$   $b_{2}=8 m$  h=7 m  $J_{x}=42x10^{6} kgm^{2}$   $J_{y}=25x10^{6} kgm^{2}$  $J_{z}=17.5x10^{6} kgm^{2}$ 

considering the structure with two vertical planes of symmetry, one longitudinal and one transversal, case for which as numerical values were proposed:

$$\begin{array}{l} a{=}7.5 \text{ m} \\ b_{3}{=}10 \text{ m} \\ b_{2}{=}10 \text{ m} \\ h{=}7 \text{ m} \\ J_{x}{=}35 x 10^{6} \text{ kgm}^{2} \\ J_{y}{=}25 x 10^{6} \text{ kgm}^{2} \\ J_{z}{=}17.5 x 10^{6} \text{ kgm}^{2} \end{array}$$

For both sets of values, the mass of the analyzed structure was considered with the value  $m = 3x10^6$  kg. Also, in both cases were proposed for study six sets of values of stiffness coefficients as follows in Table 1.

After completing the mathematical calculus, we obtain values for the parameters initially proposed - the eigenvalues and eigenvectors of the analyzed system. These values are summarized in Tables 2, 3, 4, 5.

	Var 1	Var 2	Var 3	Var 4	Var 5	Var 6
k <sub>x</sub> [N/m]	$2x10^{6}$	$4x10^{6}$	$8x10^{6}$	$16 \times 10^{6}$	$32 \times 10^{6}$	64x10 <sup>6</sup>
$k_y[N/m]$	$2x10^{6}$	$4x10^{6}$	$8x10^{6}$	$16 \times 10^{6}$	$32 \times 10^{6}$	64x10 <sup>6</sup>
$k_z[N/m]$	$8x10^{6}$	$16 \times 10^{6}$	$32 \times 10^{6}$	$64 \times 10^{6}$	$128 \times 10^{6}$	$256 \times 10^{6}$

	Case with a longitudinal vertical plane of symmetry - subsystem I									
	eigen	nat puls	freq	eigen	eigen	eigen	eigen	eigen	eigen	
	val	р	f	vect 1	vect 2	vect 3	vect 1	vect 2	vect 3	
	p <sup>2</sup>			$\mu_1$	$\mu_2$	$\mu_3$	norm	norm	norm	
							$\mu_{1n}$	$\mu_{2n}$	$\mu_{3n}$	
	89.3	9.4484	1.5038	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V1	2.4	1.5376	0.2447	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	10.3	3.2044	0.5100	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	178.5	13.3621	2.1266	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V2	4.7	2.1745	0.3461	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	20.5	4.5317	0.7212	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	357.1	18.8968	3.0075	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V3	9.5	3.0752	0.4894	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	41.1	6.4088	1.0200	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	714.2	26.7241	4.2533	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V4	18.9	4.3490	0.6922	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	82.1	9.0634	1.4425	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	1428.4	37.7936	6.0150	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V5	37.8	6.1504	0.9789	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	164.3	12.8176	2.0400	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	
	2856.7	53.4483	8.5066	0.2037	0.9990	0.0458	1.0000	1.0000	1.0000	
V6	75.7	8.6980	1.3843	0.2564	0.0416	-0.9988	1.2592	0.0416	-21.8008	
	328.6	18.1268	2.8850	0.9449	-0.0162	0.0187	4.6396	-0.0162	0.4072	

Table 2 The system parameters assessment for the subsystem I in the case with one vertical plane of symmetry

Table 3	3 The system	parameters	assessment	for the sub	osystem II 1	in the case	with one	vertical p	plane of	symmetry

Case with a longitudinal vertical plane of symmetry - subsystem II									
	eigen	nat puls	freq	eigen	eigen	eigen	eigen	eigen	eigen
	val	р	f	vect 1	vect 2	vect 3	vect 1	vect 2	vect 3
	$p^2$			$\mu_1$	$\mu_2$	$\mu_3$	norm	norm	norm
							$\mu_{1n}$	$\mu_{2n}$	$\mu_{3n}$
	2.1	1.4601	0.2324	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V1	90.0	9.4861	1.5098	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	71.5	8.4550	1.3457	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837
	4.3	2.0649	0.3286	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V2	180.0	13.4153	2.1351	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	143.0	11.9571	1.9030	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837
	8.5	2.9202	0.4648	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V3	359.9	18.9721	3.0195	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	285.9	16.9100	2.6913	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837
	17.1	4.1298	0.6573	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V4	719.9	26.8306	4.2702	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	571.9	23.9143	3.8061	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837
	34.1	5.8405	0.9295	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V5	1439.8	37.9442	6.0390	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	1143.8	33.8199	5.3826	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837
	68.2	8.2597	1.3146	-0.9996	-0.2160	-0.0023	1.0000	1.0000	1.0000
V6	2879.5	53.6612	8.5405	-0.0256	0.9079	-0.2669	0.0256	-4.2028	116.9657
	2287.6	47.8286	7.6122	-0.0105	0.3592	0.9637	0.0105	-1.6626	-422.2837

	Case with	h two vertica	l planes of sy	mmetry (longitud	inal and vertical)	subsystem I
	eigenval	nat puls	freq	eigenvect 1	eigenvect 2	eigenvect 3
	$p^2$	р	f	$\mu_1$	$\mu_2$	$\mu_3$
	2.4	1.5391	0.2450	-0.9999	-0.1830	0
V1	102.9	10.1453	1.6147	0	0	1.0000
	10.7	3.2660	0.5198	0.0160	-0.9831	0
	4.7	2.1766	0.3464	-0.9999	-0.1830	0
V2	205.9	14.3476	2.2835	0	0	1.0000
	21.3	4.6188	0.7351	0.0160	-0.9831	0
	9.5	3.0782	0.4899	-0.9999	-0.1830	0
V3	411.7	20.2905	3.2293	0	0	1.0000
	42.7	6.5320	1.0396	0.0160	-0.9831	0
	19.0	4.3532	0.6928	-0.9999	-0.1830	0
V4	823.4	28.6952	4.5670	0	0	1.0000
	85.3	9.2376	1.4702	0.0160	-0.9831	0
	37.9	6.1563	0.9798	-0.9999	-0.1830	0
V5	1646.8	40.5811	6.4587	0	0	1.0000
	170.7	13.0639	2.0792	0.0160	-0.9831	0
	75.8	8.7064	1.3857	-0.9999	-0.1830	0
V6	3293.6	57.3903	9.1340	0	0	1.0000
ľ	341.3	18.4752	2.9404	0.0160	-0.9831	0

Table 4 The system parameters assessment for the subsystem I in the case with two vertical planes of symmetry

Table 5 The system parameters assessment for the subsystem II in the case with two vertical planes of symmetry

Case with two vertical planes of symmetry (longitudinal and vertical) - subsystem II									
	eigenval	nat puls	freq	eigenvect 1	eigenvect 2	eigenvect 3			
	$p^2$	р	f	$\mu_1$	$\mu_2$	$\mu_3$			
	2.2	1.4757	0.2349	-0.9997	0.2133	0			
V1	88.2	9.3898	1.4944	-0.0262	-0.9770	0			
	71.4	8.4515	1.3451	0	0	1.0000			
	4.4	2.0869	0.3321	-0.9997	0.2133	0			
V2	176.3	13.2792	2.1135	-0.0262	-0.9770	0			
	142.9	11.9523	1.9023	0	0	1.0000			
	8.7	2.9514	0.4697	-0.9997	0.2133	0			
V3	352.7	18.7797	2.9889	-0.0262	-0.9770	0			
	285.7	16.9031	2.6902	0	0	1.0000			
	17.4	4.1739	0.6643	-0.9997	0.2133	0			
V4	705.4	26.5585	4.2269	-0.0262	-0.9770	0			
	571.4	23.9046	3.8045	0	0	1.0000			
	34.8	5.9027	0.9394	-0.9997	0.2133	0			
V5	1410.7	37.5593	5.9778	-0.0262	-0.9770	0			
	1142.9	33.8062	5.3804	0	0	1.0000			
	69.7	8.3477	1.3286	-0.9997	0.2133	0			
V6	2821.4	53.1169	8.4538	-0.0262	-0.9770	0			
	2285.7	47.8091	7.6091	0	0	1.0000			

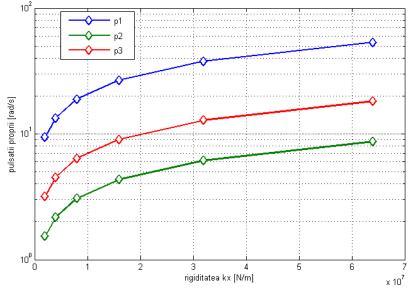


Fig. 2. The dependence between the natural pulsations of the system and the rigidity in the horizontal direction

Based on the values obtained for the three natural pulsations of the system in the six cases proposed, it was realized a graphic (Figure 2) of the dependence between the natural pulsations and the rigidity in the horizontal direction, denoted by  $k_x$ .

### 4. CONCLUSIONS

It should be mentioned that the stiffness in the x direction is equal to that in the y direction and stiffness in the z direction is a linear combi-nation of the two others. Therefore we obtained for the eigenvectors identical values, regardless of the values of the stiffness coefficient considered in calculus.

As an independent variable for the representation of the pulsation evolution was chosen the stiffness in the x direction, denoted by  $k_{x}$ .

The evolution of the pulsations corresponding to the eigenvalues follows the natural tendency imposed by the pairs of values considered for rigidities. The correlative analysis of each set of eigenvalues induces the following conclusion, namely that linear combinations between the stiffness in the horizontal plane and the one in the vertical plane require similar evolutions.

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